

A Parametric Comparison Between the Behavior of $1S_0$ and $3S_1$ of Helium and the Excited S-States of Helium Neon laser



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Abstract

Gas-discharges have a wide range of spectroscopic parameters including the after-glow excitations. If a fixed mixture of He and Ne is operated as a traditional discharge phenomenon, then the important parameters could be measured spectroscopically. This project aims to parametric comparison between the situation of both metastables of He states and the excited S-states of Ne in terms of different collision cross-sections, life-time considerations, spectral notations and many other resonant transfer parameters. The results were expected to be favorable and suitable for our purposes and they were optimum for a mixture of He and Ne in low power gas lasers.

Keyword : Meta-stable states, spectroscopy of inert gases, laser physics

Introduction

The He-Ne lasers active medium is an excellent example to be studied on the behavior of the process of exciting the $1S_0, 3S_1$ of helium and 2S, 3S of neon.

All inert gases, He, Ne, Ar, Kr, and Xe possess at least one metastable state for which no selection rule, in both J-J and L-S coupling schemes, will be obeyed. The reason of this is that the atoms or molecules situated in these levels stay as much as they can and after this period they will de-excite to the lower state either by collision with the container wall or with neighboring atom. The optimum life time for these states range from few micro seconds to ~ 10 ms. Life time consideration will be of great importance in this work and the tendency will be toward the mechanism of excitation transfer at resonance, i.e. the net number of absorptions per unit time is the same as the net number of emissions per unit time also.

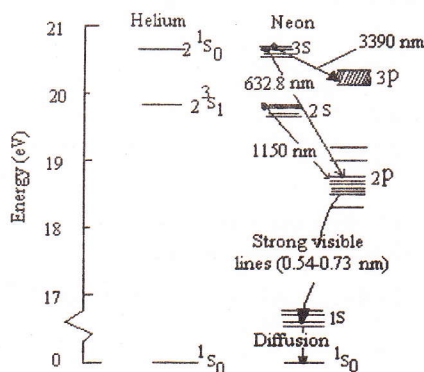
The factors responsible for this resonance were studied in detail and theoretical

calculations were done to compare them with experimental one. The exopton of a

mixture of helium and neon as the active medium in gas laser was analyzed to identify ideal requisites to operate this laser.

Theory

Gas lasers [1] were generally employed a mixture of non- reacting gases to act as an active medium with specified ratios, for instance in CO₂ lasers the ratio of CO₂ : N₂ : He is 1:2:3



in volume proportion.

Partial atomic energy level diagram of helium and neon is shown in fig.1

In fig.(1) both 1S_0 and 3S_1 in helium are metastables there life time estimated to be about 8 μs and 125 μs respectively, So the chance of 1S_0 to collide with neighboring atoms is much less than that of 3S_1 , that stays for much longer time rather than 1S_0 . The metastable of neon is represented by the 1s state where there is no lower states below it, thus the atoms in this state cool down through collision. Spectroscopic differences between metastables of helium and (s) states of neon are shown in fig.(2) In this figure, the spectroscopic equations governing the excitation transfer between the levels are reduced by two reactions :

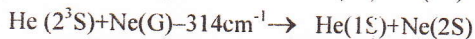
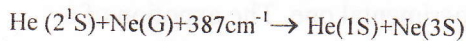
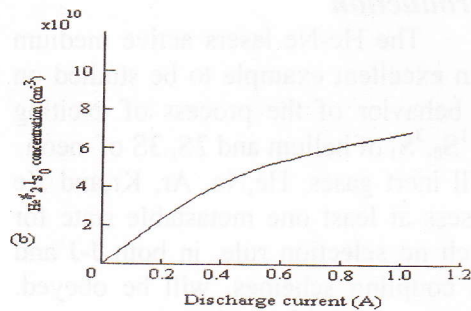
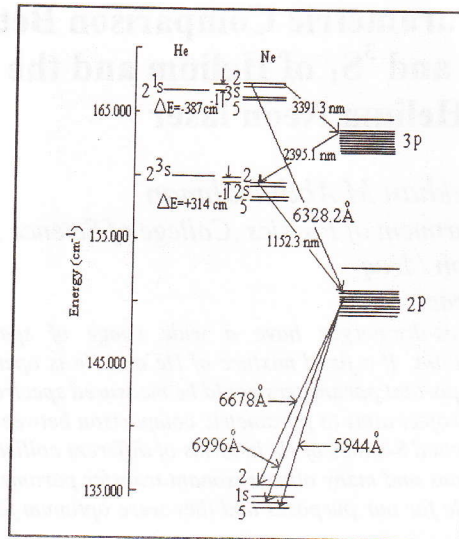


Fig.(2) Spectroscopic details between meta-stables of helium and (s) states of neon This means that 3S state of neon lays above 1S_0 of helium by $387 \text{cm}^{-1} \sim (0.048 \text{ eV})$ and 2S below 3S_1 of He by $314 \text{cm}^{-1} (0.039 \text{ eV})$.

The collision excitation transfer in such mixtures are very complex we shall focus only on those which occur at resonance.



Fig(2) spectroscopic details between meta-stables of helium and (s) states of neon.

Transfer mechanisms

Verdyen et al [2] have suggested that the number of variables in such transfers should be reduced by four processes. Fig.3 shows detail of these resonance transfers between the metastable ($^1S_0, ^3S_1$) of helium and ($3s, 2s$) of neon.

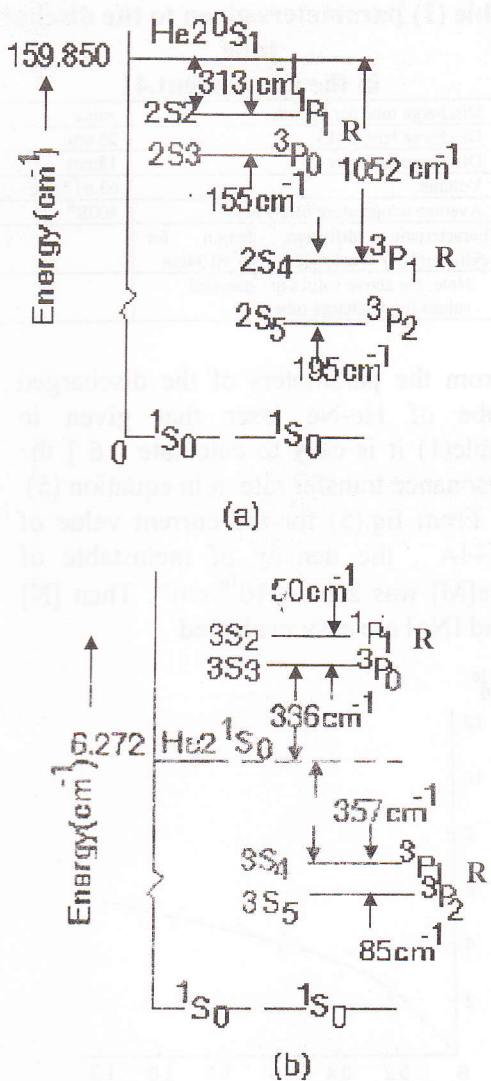


Fig.(3) The resonance transfer levels between :a- 3S_1 and $2s$. b- 1S_0 and $3s$
 * R : refer to resonance state

Investigations by Phelps and his colleagues [3] shows that for a simplified model of He-Ne laser the current of the discharge for all levels varies with densities of states semi-exponentially, i.e. for N_0, N_1, N_2 , which are the metastables $1S, 2S$ and $3S$ of Ne respectively and M is the meta-stable of He (N_0, N_1, N_2 and M are density numbers of excited atoms). The main equations for these states are :

$$d[M]/dt = n_e v_e [He] - M/\tau_m - r_i[M][Ne] - n_e v_e [M] \dots (1)$$

$$d[N_2]/dt = r_i[M][Ne] - A_2[N_2] \dots (2)$$

$$d[N_1]/dt = n_e v_e [Ne] + n_e i_0 v_e [N_0] - A_1[N_1] \dots (3)$$

$$d[N_0]/dt = n_e v_e [Ne] + A_1[N_1] - n_e v_e [N_0] + A_2[N_2] - [N_0]/\tau_0 - n_e v_e [N_0] \dots (4)$$

Here: 1-All quantities in brackets refer to densities (number of excited atoms per unit volume) where:

- v_e - the electron velocity calculated from $v_e = \sqrt{2T/m}$ where T is the kinetic energy of the electron and m is its mass.
- n_e - the density of the state.
- σ - the cross-section of the transition.
- τ - the state life time
- r_i - the transition rate.

2-The products of ()s i.e. (n_e, i, σ, \dots) , (N) and (n_e) represented the collision frequency of any process and multiplied by the proper density of the state.

Jeffery and Dalgarno [4], have shown algebraic analysis for these equations and used the Rapnsden-Newton iteration of the following values for the density of states, they were obtained by :

$$M = n_{em} \text{Ne}[\text{He}] / \tau_m^{-1} + r_i[\text{Ne}] + n_{ei} \text{Ne} \dots (5)$$

$$N_0 = n_e [n_0 \text{Ne} + n_1 \text{Ne}] [\text{Ne}] / \tau_0^{-1} + n_e n_3 \text{Ne} \dots (6)$$

$$N_1 = n_e n_1 \text{Ne}[\text{Ne}] + n_e n_{01} \text{Ne}[\text{Ne}] \dots (7)$$

$$N_2 = r_i[\text{Ne}][M] / \tau_2^{-1} \dots (8)$$

If these equations hold for an optimum discharge current and voltage then we would be in a position to calculate the transfer rate (r_i) [5] for both absorption and emission transfers.

Experimentaldetails

The apparatus arranged in fig.(4) was used to obtain the experimental part of this work. A He-Ne laser of 20mW [out put] connected with two digital multi-meters one in series and the other in parallel with the output of the step-up transformer of the laser cavity

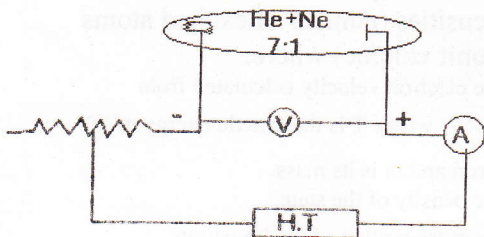


Fig.(4) Experimental arrangement of the He-Ne discharge tube

When the step-up was established against any mechanical shock, then through the rheostat the current was varied through 0.05A and the voltmeters reading was observed.

A repetitive set of readings were taken then after lengthly observation for the optimum temperature of the gas discharges

the current was stabilized at 0.44A and at potential difference of 1100V the power consumption obtained was.

$$P = V * I = 1100 * 0.44 = 484 \text{ Watts} \dots (9)$$

If the volume of the discharge is known then the density of metastables of He [M] and (No) of Ne are calculated.

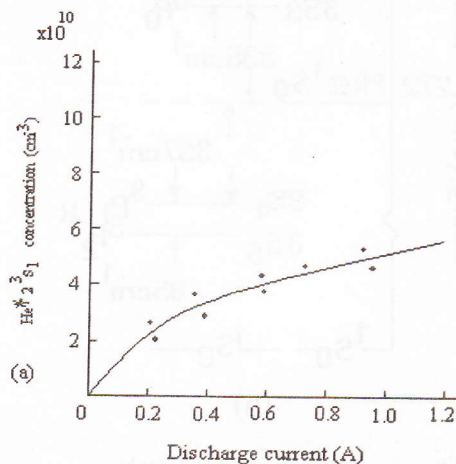
The results and tabulated in table No 2

Table (1) parameters given to the discharge tube in the experiment.41

Discharge tube parameters	value
Discharge length (L)	25 cm.
Discharge diameter (2r)	18mm.
Volume	63.617 cm ³
Average temperature of the tube	400K ⁰
Characteristic diffusion length for cylinder $(1/\lambda)^2 = (\pi/4)^2 + (2.405/r)^2 = 0.34 \text{cm}$	
Note :the above values are standard values for discharge tube	

From the parameters of the discharged tube of He-Ne laser that given in table(1) it is easy to calculate [6] the resonance transfer rate r_i in equation (5)

From fig.(5) for the current value of 0.44A , the density of metastable of He[M] was $2.11 \times 10^{16} \text{ cm}^{-3}$. Then [N] and [No] are easy evaluated



State electron densities of metastables with discharge current.

Results and Discussion

From equations (5) to (8) the discharge current is stabilized, then the relation between $[Ne]$ and the decay rate $r_t \times 10^6 \text{sec}^{-1}$ is plotted as shown in fig.(6)

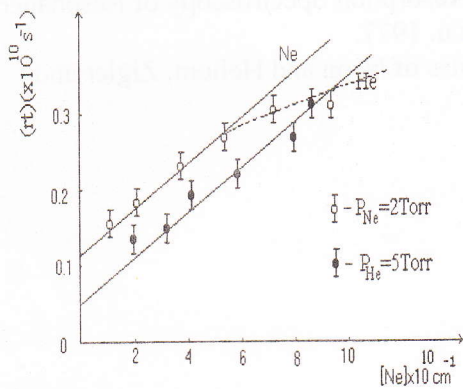


Fig.(6) Decay rate of meta-stable rate (r_t) with density of Ne

It is clear from fig.(6) that for low $[Ne]$ densities the transfer rate between metastables of He and 3S of Ne is low, and increases exponentially at about $6.85 \times 10^{10} \text{cm}^{-3}$ for $[Ne]$. Then we calculate the rate coefficient of He at resonance for both metastable states

Once Boltzmann distribution [6] is used to calculate the atomic densities for 2S and 3S in Ne without discharge by using eq.(9) and the standard values of the discharge tube in table(1) the ionization cross-section for both states are used to find the electron velocities at 1S_0 and 3S_1 for He. The results are arranged in table (2).

Table(2) Theoretical calculation for collision process between $1S_0 \rightarrow 3S$ and $3S_1 \rightarrow 2S$ in resonance tube

Parameter	1S_0	3S_1
σ_i (cm^2) ionization cross section	0.28×10^{-16}	4.1×10^{-18}
v_e cm s^{-1} the electron velocity	2.639×10^8	2.692×10^8
n_e cm^{-3} the density of state	4.1×10^{13}	0.42×10^{11}
r_e emission rate	$3.4 \times 10^{11} \text{cm}^{-3} \text{s}^{-1}$	$3.7 \times 10^{11} \text{cm}^{-3} \text{s}^{-1}$
r_a absorption rate	$6.46 \times 10^{11} \text{cm}^{-3} \text{s}^{-1}$	$3.85 \times 10^{11} \text{cm}^{-3} \text{s}^{-1}$

Conclusion

From the result of calculations it is obvious that there are few spectroscopic differences between the two metastables of He. The important conclusion regarding these states are :

1. The total number of collision per unit time [collisional frequencies] in 1S_0 is much less than that of 3S_1 , the reason for this is that the latter has a better chance to collide with atoms of 2S in Ne rather than with 3S

2. The absorption and emission transition rates $[r_a, r_e]$ are characteristics of collision cross-section at resonance between meta-stable states of [He] and (S) states of Ne.

This means that they are independent of the discharge parameters like tube length, cross-section, temperature or densities but they are functions of excitation and ionization cross-sections only.

3. In this project the important parameters of the discharge can not be measured classically. But when the equations of state (eq.1 \rightarrow eq.4) are used then the microscopic views of the discharge could be studied in terms of the concentrations.

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بهراوردى نه نيون خوماليه كانى ناستى S_0^1 و S_1^3 ي هيليووم و ناستى بزوينراوى كبهراوردى نيون نه له يزهري هيليووم نيون دا .

بهري خان محمد مهدي عبدالرحمن

فيزيك / كوليجي زانست / زانكوي سليمانى / هه ريهي كوردستان - عيراق

پوخته

ناشكرايه كه كردارى خالى كردنه وهى كارها نه گاذا كارى گهريكى فراوانى ههيه وگورانكاريهكى زور نه دوايى خاليكردنه وه كه دا رو نه دا. وهك بزواندنى ناوهنده كه وگوراني ناستى ووزه. نه گهر چه سپهر ريزه به كه نه هيليووم و نيون تيكه ل بگرين نه نيو تيوييكي خاييكه ره وهى كاره بايى دا نه وا شه بهنگ كارى تيشكى به دهست هينراو زانباريهكى زورمان نه داتى. مه بهست نه م تويزينه وهيه بهراورد كردنى رولى ههردوو ناستى نينچه جيگري هيليووم و ناسته وروژا وه كانى جورى S نه گاذا نيون دا به به هاوكارى چونه تى گونجاني روويهري بهريه كه كه وتن, نه گهري بهريه كه كه وتن, ته مه نى ناسته كان, هيمايى شه به نكيه كان وگه لى روانه ت كارى ترى زرنكه گواستنه وه. نه نجامه كان زور گونجاو بوون نه گه ل هاوتاتيو ريه كانيان و توانرا نرخی نه باربان بو بدوزرينه وه نه له يزهري هيليووم نيونى توانانز مه كاندا.

الخلاصة

دراسة مقارنة لسلوك المستويين والمستوي S للنيون في ليزر الهيليوم_ نيون

S_0^1 و S_1^3 للهيليوم

به ري خان محمد عبدالرحمن

قسم الفيزياء / كلية العلوم / جامعة السليمانية / اقليم كردستان - العراق

تعد عملية التفريغ الغازي وسيلة ملائمة لإزالة الأوساخ الغازية، فهناك تغيرات كثيرة تطرأ على الوسط الغازي بعد عملية التفريغ من ضمنها تغير مستوي طاقة مكونات الوسط الغازي. عند دمج نسبة محددة من غازي الهيليوم والنيون في انبوبة التفريغ الغازية فان التعليل الطيفي للضوء المنبعث يعطي معلومات قيمة حول الموضوع. الهدف من هذه الدراسة هي مقارنة المستويين الشبه المستقرين للهيليوم مع المستويات المثارة من نوع S للنيون بدلالة المقطع العرضي للتصادم، مدى العمر، الرموز الطيفية وغيرها من بارامترات الانتقال الرنيني. نتائج الدراسة كانت مرضية ومناسبة لأغراضنا ومثالية في حالة ليزرات الهيليوم_ نيون الغازية ذات القدرات الواطنة.